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### Abstract

The oxides  $Ln_3Li_5Sb_2O_{12}$  ( $Ln \equiv Pr$ , Nd, Sm) were prepared in air by heating a mixture of  $Ln_2O_3$ ,  $LiNO_3$  and  $Sb_2O_3$  in the temperature range 1023-1243 K. Lattice parameters as well as atomic coordinates in the space group  $I2_13$  (Z=8) are established from Xray powder diffraction data by the Rietveld method. Magnetic susceptibilities from 4.2 to 300 K follow a Curie–Weiss law above 50 K ( $Ln \equiv Pr$ ) or 70 K ( $Ln \equiv Nd$ ). This is attributed to the splitting of the ground state associated with the  $Ln^{3+}$  ions by the influence of the crystal field. The magnetic moments, 3.54 and 3.61  $\mu_B$  for praseodymium and neodymium respectively, agree with those calculated by Hund's formula. The Sm<sup>3+</sup> behaviour can be explained by taking into account that the splitting of the multiplets is not too large compared to kT. The magnetic moment observed at room temperature for this cation is 1.6  $\mu_B$ .

# 1. Introduction

A new series of mixed oxides,  $Ln_3Li_5Sb_2O_{12}$  ( $Ln \equiv La$ , Pr, Nd, Sm, Eu), has recently been synthesized [1]. The unit cell is cubic (space group  $I2_13$ , Z = 8).

In this paper we report the experimental magnetic susceptibilities for the oxides  $Ln_3Li_5Sb_2O_{12}$  ( $Ln \equiv Pr$ , Nd, Sm) in polycrystalline samples, which were also characterized by X-ray powder diffraction data, refined on the basis of the Rietveld analysis method.

### 2. Experimental details

 $Ln_3Li_5Sb_2O_{12}$  ( $Ln \equiv Pr$ , Nd, Sm) samples were prepared from  $Ln_2O_3$ ,  $LiNO_3$ and  $Sb_2O_3$  (Merck). The initial mixture in stoichiometric ratio is heated to 1023 K for 6 h. The second step is a thermal treatment at 1243 K. In this process, sublimation of lithium oxide takes place and it is necessary to add an excess of 30%–40% LiNO<sub>3</sub>. Therefore the stoichiometry of this compound could not be fixed, but the tentative chemical formula of the phase  $La_3Li_5Sb_2O_{12}$  [2] is proposed taking into account the structural analogy with the niobium (V) and tantalum (V) ternary oxides [3].

## 3. Results and discussion

Powder X-ray diffraction patterns were registered at a rate of 0.1° (2 $\theta$ ) per minute by means of a Siemens Kristalloflex diffractometer powered by a D500 generator using nickel-filtered Cu K $\alpha$  radiation. A 2 $\theta$  step scan of 0.04° was used and the Rietveld profile analysis method [4] was employed for refinement of X-ray diffraction patterns in the range 10°-120° (2 $\theta$ ) for 287 observed reflections. The structure of La<sub>3</sub>Li<sub>5</sub>(Nb, Ta)<sub>2</sub>O<sub>12</sub> was used as a trial model. The Rietveld programme minimizes the function  $\chi^2 = (R_{wP}/R_{EXP})^2$ . The reflection conditions are h+k+l=2n for (h, k, l), k+l=2n for (0, k, l) and h=2n for (h, 0, 0), compatible with the space group  $I2_13$ . The structural data for Ln<sub>3</sub>Li<sub>5</sub>Sb<sub>2</sub>O<sub>12</sub> (Ln = Pr, Nd, Sm) are given in Table 1. The unit cell parameters decrease linearly in accordance with the lanthanide contraction effect. Table 2 lists the experimental bond lengths and angles of Ln<sub>3</sub>Li<sub>5</sub>Sb<sub>2</sub>O<sub>12</sub> (Ln = Pr, Nd, Sm).

The structure of these compounds is isomorphous to  $La_3Li_5Sb_2O_{12}$  and can be described as a three-dimensional framework of lanthanide, antimony and oxygen atoms. The SbO<sub>6</sub> octahedra are located along the threefold axes of the unit cell and are associated with three lanthanide atoms, forming infinite chains of composition  $(Ln_3Sb_2O_{12})_n$ . These chains are bonded with others by means of lanthanide and lithium atoms. Figure 1 shows the orientation of  $(Ln_3Sb_2O_{12})_n$  chains along a threefold axis.

Magnetic susceptibility measurements were made using the Faraday method in the temperature range 4.2–300 K with a DMS5 Pendule susceptometer. The maximum magnetic field applied was 1 kG with H(dH/dZ) = 29 kG<sup>2</sup> cm<sup>-1</sup>. The equipment was calibrated with Hg(Co(SCN)<sub>4</sub>) and Cd<sub>2</sub>(SO<sub>4</sub>)<sub>3</sub>·8H<sub>2</sub>O and  $\chi$  was independent of the field in the temperature range of the measurements. Values of effective moment ( $\mu_B$ ) and Weiss constant ( $\theta$ ) were obtained by least-squares fits from the linear part of the  $\chi^{-1}$  vs. T plots.

The magnetic susceptibility of  $Pr_3Li_5Sb_2O_{12}$  follows a Curie–Weiss law behaviour from 300 to 50 K. Below this temperature deviations upwards from linearity are observed, analogous to those found in some compounds of praseodymium [5, 6]. Since this cation has an even number of f electrons, the lowest crystal field level could be a singlet and the upward deviations observed at the lower temperatures are probably due to the population of this singlet level (see Fig. 2). The reciprocal molar magnetic susceptibility for Nd<sub>3</sub>Li<sub>5</sub>Sb<sub>2</sub>O<sub>12</sub> obeys a Curie–Weiss law,  $\chi = 1.63/T + 14.8$ , above 70 K and the moment calculated from the  $\chi^{-1}$  vs. T plot suggests that crystal fields effects are probably present (Fig. 2). Similar remarks have previously been reported for different neodymium compounds [7–9].

The observed and calculated magnetic moments and Weiss constants for  $Ln^{3+}$  in the oxides can be found in Table 3. In these cases negative

Ln ≡ Pı					Ln = Nc					Ln≡Sm				
Atom	Site	x	ų	ĸ	Atom	Site	x	y	Ń	Atom	Site	x	y	×
Pr(1)	12b	0.118(9)	0.000	0.250	(1)pN	12b	0.124(2)	0.000	0.250	Sm(1)	12b	0.116(6)	0.000	0.250
Pr(2)	12b	0.619(8)	0.000	0.250	Nd(2)	12b	0.622(9)	0.000	0.250	Nd(2)	12b	0.620(9)	0.000	0.250
Sb(1)	8a	0.000	0.000	0.000	Sb(1)	8a	0.000	0.000	0.000	Sb(1)	8a	0.000	0.000	0.000
Sb(2)	8a	0.250	0.250	0.250	Sb(2)	8a	0.250	0.250	0.250	Sb(2)	8a	0.250	0.250	0.250
Li(1)	24c	0.630(0)	0.307(2)	0.533(3)	Li(1)	24c	0.603(6)	0.307(4)	0.533(1)	Li(1)	24c	0.534(1)	0.284(2)	0.612(9)
Li(2)	24c	0.454(4)	0.383(3)	0.699(9)	Li(2)	24c	0.454(9)	0.382(5)	0.696(7)	Li(2)	24c	0.486(9)	0.385(1)	0.632(7)
0(1)	24c	0.323(0)	0.088(1)	0.241(4)	0(1)	24c	0.263(0)	0.094(8)	0.204(5)	0(1)	24c	0.310(8)	0.101(7)	0.225(1)
0(2)	24c	0.218(4)	0.395(4)	0.297(8)	0(2)	24c	0.234(4)	0.396(4)	0.332(9)	0(2)	24c	0.219(7)	0.423(0)	0.284(1)
0(3)	24c	0.312(0)	0.544(4)	0.430(4)	0(3)	24c	0.378(6)	0.544(2)	0.453(6)	0(3)	24c	0.360(8)	0.514(1)	0.423(6)
0(4)	24c	0.614(2)	0.470(5)	0.561(0)	0(4)	24c	0.663(2)	0.438(7)	0.529(2)	0(4)	24c	0.633(2)	0.460(8)	0.575(3)
Cell pi Volum Reliabi R <sub>P</sub> = 19	tramet <sup>i</sup> $({\rm Å}^3)$ : lity fac .6, $R_{\rm W}$	er (Å): $a = 2070.23(4)$ to 2070.23(4) to 2000; $R_1$	12.745(2) t) $B = 8.33, R_{\rm F}$	.= 7.86	Cell pa Volume Reliabi $R_{\rm P}=18$	tramet e ( $\mathring{A}^3$ ): lity fa	cr (Å): $a = 2035.82(1)$ zors (%): a = 21.0, R	12.674(5) .) $_{\rm B} = 6.22, R_{\rm I}$	·=6.67	Cell pa: Volume Reliabil $R_p = 18.$	ramete $(Å^3)$ : ity fac	r (Å): <i>a</i> = 1 1989.44(1) tors (%): = 20.0, <i>R</i> <sub>B</sub>	2.577(0) = 8.87, $R_{\rm F}$	= 7.32

Structural parameters for Ln<sub>3</sub>Li<sub>5</sub>Sb<sub>2</sub>O<sub>12</sub>

TABLE 1

TABLE 2				
Bond lengths (angs	troms) and b	ond angles (	(degrees) of Ln <sub>3</sub> Li <sub>5</sub> Sb <sub>2</sub> C	12
Ln≡ Pr		***		
Sb(1)O(3)	2.61(7)	3×	Pr(1)O(4)	2.43(9)
Sb(1)O(4)	1.69(1)	$3 \times$	Pr(1) - O(4)	2.60(6)
Sb(2)-O(1)	2.26(7)	$3 \times$	Pr(1) - O(1)	2.83(3)
Sb(2)-O(2)	1.90(5)	$3 \times$	Pr(1) - O(1)	2.90(4)
Pr(2)-O(2)	2.62(5)	$2 \times$	Pr(2) - O(2)	2.54(2)
Pr(2)-O(3)	2.39(5)	$2 \times$	Pr(2)-O(3)	2.52(0)
Li(1) - O(1)	2.01(0)		Li(2)-O(1)	2.02(2)
Li(1) - O(2)	2.19(5)		Li(2)-O(1)	2.34(3)
Li(1) - O(2)	2.69(5)		Li(2)-O(2)	2.19(8)
Li(1)-0(3)	2.21(1)		Li(2)-O(3)	2.19(4)
Li(1)-O(3)	2.00(0)		Li(2) - O(4)	2.49(6)
Li(1)-O(4)	2.11(2)		Li(2)–O(4)	2.52(2)
O(4)-Pr(1)-O(4)		90.02	O(1)-Pr(1)-O(1)	
O(2) - Pr(2) - O(2)		70.43	O(3) - Pr(2) - O(3)	
O(3)-Sb(1)-O(4)		84.48	O(4)-Sb(1)-O(3)	
O(1)-Sb(2)-O(2)		89.41	O(2)-Sb(2)-O(1)	
O(3)-Li(1)-O(1)		65.67	O(1)-Li(1)-O(3)	
O(3)-Li(1)-O(2)		94.95	O(2)-Li(1)-O(2)	
O(2)-Li(1)-O(4)		93.04	O(4)-Li(1)-O(3)	
O(2)-Li(2)-O(4)		92.26	O(4)-Li(2)-O(1)	
O(1)-Li(2)-O(1)		63.08	O(1)-Li(2)-O(3)	
O(3)-Li(2)-O(4)		89.74	O(4)-Li(2)-O(2)	
Ln≡Nd				

O(1)–Li(2)–O(1) O(3)–Li(2)–O(4)		$63.08 \\ 89.74$	O(1)-Li(2)-O(3) O(4)-Li(2)-O(2)		55.77 90.02
Ln≅Nd		_			
Sb(1)-O(3)	1.73(9)	3×	Nd(1)O(4)	2.62(2)	$2 \times$
Sb(1)-O(4)	2.24(0)	$3 \times$	Nd(1)O(4)	2.94(6)	$2 \times$
Sb(2)-O(1)	2.06(0)	$3 \times$	Nd(1)O(1)	2.24(0)	$2 \times$
Sb(2)-O(2)	2.14(1)	$3 \times$	Nd(1)-O(1)	2.62(6)	$2 \times$
Nd(2)O(2)	2.14(1)	$2 \times$	Nd(2)-O(2)	2.46(9)	$2 \times$
Nd(2)O(3)	2.64(0)	$2 \times$	Nd(2)-O(3)	2.73(6)	$2 \times$
Li(1) - O(1)	2.55(1)		Li(2)-O(1)	2.14(6)	
Li(1) - O(2)	2.79(9)		Li(2)-O(1)	2.38(6)	
Li(1)-O(2)	2.09(6)		Li(2)-O(2)	2.42(7)	
Li(1)-O(3)	2.10(4)		Li(2)-O(3)	2.09(6)	
Li(1)-O(3)	2.46(3)		Li(2)-O(4)	2.08(0)	
Li(1)–O(4)	2.20(1)		Li(2)O(4)	2.01(7)	
O(4)-Nd(1)-O(4)		78.47	O(1)-Nd(1)-O(1)		68.29
O(2)-Nd(2)-O(2)		72.06	O(3)-Nd(2)-O(3)		84.57
O(3)-Sb(1)-O(4)		77.65	O(4)-Sb(1)-O(3)		89.61
O(1)-Sb(2)-O(2)		92.79	O(2)-Sb(2)-O(1)		69.99
O(3)-Li(1)-O(1)		91.05	O(1)-Li(1)-O(3)		79.74
O(3)-Li(1)- $O(2)$		110.90	O(2)-Li(1)-O(2)		89.02
O(2)-Li(1)-O(4)		78.26	O(4)-Li(1)-O(3)		64.27
O(2)-Li(2)-O(4)		89.80	O(4)-Li(2)-O(1)		108.90
O(1)-Li(2)-O(1)		92.81	O(1)-Li(2)-O(3)		85.78
O(3)-Li(2)-O(4)		88.97	O(4)-Li(2)-O(2)		90.20

(continued)

 $2 \times$ 

 $140.80 \\ 59.69 \\ 87.54 \\ 110.38 \\ 81.44 \\ 73.68 \\ 91.55 \\ 124.03$ 

TABLE 2	. (continued)
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$Ln \equiv Sm$					
Sb(1)=O(3)	2.00(3)	3×	Sm(1)-O(4)	2.62(5)	$2 \times$
Sb(1)-O(4)	1.90(7)	$3 \times$	Sm(1)-O(4)	2.26(1)	$2 \times$
Sb(2)-O(1)	2.04(1)	$3 \times$	Sm(1)-O(1)	2.77(4)	$2 \times$
Sb(2)-O(2)	2.24(0)	$3 \times$	Sm(1)-O(1)	2.93(8)	$2 \times$
Sm(2)-O(2)	2.26(6)	$2 \times$	Sm(2)-O(2)	2.79(7)	$2 \times$
Sm(2)O(3)	2.20(2)	$2 \times$	Sm(2)-O(3)	2.39(8)	$2 \times$
Li(1)-O(1)	2.09(4)		Li(2)O(1)	2.12(3)	
Li(1)-O(2)	2.84(2)		Li(2)O(1)	2.40(6)	
Li(1)-O(2)	2.10(0)		Li(2)-O(2)	2.19(3)	
Li(1)-O(3)	2.09(9)		Li(2)-O(3)	2.09(6)	
Li(1)-O(3)	2.27(1)		Li(2)-O(4)	2.41(4)	
Li(1)-O(4)	2.39(1)		Li(2)-O(4)	2.19(5)	
O(4)Sm(1)O(4)		66.72	O(1)-Sm(1)-O(1)		69.03
O(2)-Sm(2)-O(2)		77.54	O(3)-Sm(2)-O(3)		61.92
O(3)-Sb(1)-O(4)		102.85	O(4)-Sb(1)-O(3)		90.38
O(1)-Sb(2)-O(2)		89.54	O(2)-Sb(2)-O(1)		77.62
O(3)-Li(1)-O(1)		88.81	O(1)-Li(1)-O(3)		83.96
O(3)-Li(1)-O(2)		98.99	O(2)-Li(1)-O(2)		79.44
O(2)-Li(1)-O(4)		99.42	O(4)-Li(1)-O(3)		84.55
O(2)-Li(2)-O(4)		73.07	O(4)-Li(2)-O(1)		137.57
O(1)-Li(2)-O(1)		70.70	O(1)-Li(2)-O(3)		119.33
O(3)-Li(2)-O(4)		96.61	O(4)-Li(2)-O(2)		88.81



Fig. 1. A threefold axis along which chains of  $(Ln_3Sb_2O_{12})_n$   $(Ln \equiv Pr, Nd, Sm)$  are disposed and location of lithium atoms (circles).

values for the Weiss constant  $(\theta)$  are found. It is obvious that the Weiss constant values are due to the crystal field effects.

A different behaviour in the temperature dependence of the magnetic properties of  $Sm_3Li_5Sb_2O_{12}$  is observed (Fig. 3). The susceptibility of



Fig. 2. Temperature dependence of reciprocal magnetic susceptibility for  $Ln_3Li_5Sb_2O_{12}$  per mole of  $Ln^{3+}$  ( $Ln \equiv Pr$ , Nd).



Fig. 3. Temperature variation of magnetic susceptibility per mole of  $Sm_3Li_5Sb_2O_{12}$ . Full lines represent calculated values.

 $Sm_3Li_5Sb_2O_{12}$  remains almost constant between 100 and 240 K. This characteristic behaviour is attributed to the fact that the energy differences of the <sup>6</sup>H multiplet levels for  $Sm^{3+}$  are comparable to kT, and for this reason Van Vleck can be used in a first approximation [10], neglecting the *J*-mixing and crystal-field-splitting effects, although this gives a poor fit of the ex-

Ln	Ground state of the free ion	$rac{\mu^{ m c}}{(\mu_{eta})}$	$rac{\mu^e}{(\mu_B)}$	<i>θ</i> (К)
Pr	<sup>3</sup> H <sub>4</sub>	3.58	3.54	- 16.0
Nd	<sup>4</sup> I <sub>9/2</sub>	3.62	3.61	-14.8
Sm	<sup>6</sup> H <sub>5/2</sub>	1.6	1.6ª	-

 TABLE 3

 Magnetic parameters for Ln<sub>3</sub>Li<sub>5</sub>Sb<sub>2</sub>O<sub>12</sub>

<sup>a</sup>Calculated at 298 K.

perimental and calculated susceptibilities, as Fig. 3 shows, and a value of  $\lambda = 298 \text{ cm}^{-1}$  for the orbit-coupling constant is obtained.

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